

1 **An improved empirical model of electron and ion fluxes at geosynchronous orbit**
2 **based on upstream solar wind conditions.**

3
4 *M. H. Denton^{1,2}, M. G. Henderson³, V. K. Jordanova³, M. F. Thomsen⁴,*
5 *J. E. Borovsky^{1,5}, J. Woodroffe³, D. P. Hartley⁶, and D. Pitchford⁷.*

6
7 1. Center for Space Plasma Physics, Space Science Institute, CO 80301, USA.

8 2. New Mexico Consortium, Los Alamos, NM 87544, USA.

9 3. ISR-1, Los Alamos National Laboratory, NM 87545, USA.

10 4. Planetary Science Institute, AZ 85719, USA.

11 5. University of Michigan, Michigan, Ann Arbor, MI 48109, USA.

12 6. University of Iowa, Iowa City, IA 52242, USA.

13 7. SES Engineering, L-6815 Château de Betzdorf, Luxembourg.

ABSTRACT

16

17 A new empirical model of the electron fluxes and ion fluxes at geosynchronous orbit (GEO) is
18 introduced, based on observations by Los Alamos National Laboratory (LANL) satellites. The model
19 provides flux predictions in the energy range ~ 1 eV to ~ 40 keV, as a function of local-time, energy,
20 and the strength of the solar-wind electric field (the negative product of the solar wind speed and the z-
21 component of the magnetic field). Given appropriate upstream solar-wind measurements, the model
22 provides a forecast of the fluxes at GEO with a ~ 1 hour lead time. Model predictions are tested against
23 in-sample observations from LANL satellites, and also against out-of-sample observations from the
24 CEASE-II detector on the AMC-12 satellite. The model does not reproduce all structure seen in the
25 observations. However, for the intervals studied here (quiet and storm times) the Normalized-Root-
26 Mean-Squared-Deviation (NRMSD) $< \sim 0.3$. It is intended that the model will improve forecasting of
27 the spacecraft environment at GEO and also provide improved boundary/input conditions for physical
28 models of the magnetosphere.

29

30 **1. Introduction**

31 Geosynchronous orbit (GEO), at a radial distance of $6.6 R_E$ (Earth radii), is one of the most popular
32 locations for communications, scientific, and military satellites (see Figure 1). This is primarily due to
33 the fact that satellites located in this orbit have an orbital period of 24 hours, allowing them to remain
34 at the same geographic longitude above the Earth during their operational lifetime. Predictions of the
35 plasma environment encountered by satellites at GEO [Purvis *et al.*, 1984; O'Brien and Lemon, 2007;
36 Thomsen *et al.*, 2007; Sicard-Piet *et al.*, 2008; O'Brien, 2009; Ginet *et al.*, 2014; Hartley *et al.*, 2014;
37 Ganushkina *et al.*, 2013; 2014; 2015; Denton *et al.*, 2015] provide spacecraft designers and operators
38 with estimates of the plasma conditions (e.g. the ion flux and the electron flux) that satellite hardware
39 will be subjected to on orbit. If such predictions are based on upstream solar-wind conditions (e.g.
40 measured by the ACE satellite or the DSCOVR satellite situated in Lissajous orbits at the L1
41 Lagrangian point between the Earth and the Sun) then this allows a lead time of around one hour from
42 the flux predictions being made to when such fluxes may be encountered. Since elevated fluxes are
43 generally considered a hazard for satellites, a lead time of around one hour can be used to potentially
44 take remedial action with the intention of mitigating damaging effects upon the satellite hardware.
45 Understanding the environment at GEO is one scientific topic where the operational community and
46 the scientific community both invest significant effort and where each communities priorities may be
47 aligned [O'Brien *et al.*, 2013].

48

49 In addition to the hardware-related uses of electron and ion flux predictions, a variety of scientific
50 models of the inner magnetosphere also use fluxes at GEO as their outer boundary conditions (e.g.
51 Jordanova *et al.* [1998; 2003], Zaharia *et al.* [2005; 2006], Katus *et al.* [2014]). Hence, development
52 of improved predictions of the fluxes at GEO has the potential to benefit both the scientific and

53 operational communities.

54

55 In a recent study we introduced a new model of the ion and electron fluxes at GEO in the energy range
56 ~ 1 eV to ~ 40 keV as a function of local time, geomagnetic activity, and solar-activity [*Denton et al.*,
57 2015]. The model is based on observations made between 1989 and 2007 by seven Los Alamos
58 National Laboratory (LANL) satellites based at GEO. Magnetospheric Plasma Analyzer (MPA)
59 instruments (electro-static analyzers) onboard the satellites measure both the electron and the ion
60 energy-per-charge distributions between ~ 1 eV/q and ~ 40 eV/q [*Bame et al.*, 1989; *Thomsen et al.*,
61 1999]. In brief, each point in the entire MPA dataset (over 80 satellite-years of data) was allocated to
62 the appropriate bin based on an array of 40 energies (equally spaced logarithmically between 1 eV and
63 40 eV), 24 local-times, and 28 discrete values of the Kp index [*Bartels et al.*, 1939; *Thomsen*, 2004],
64 for both ions and electrons. Solar activity variations were included in the model by carrying out the
65 above binning for four ranges of the F10.7 index (all F10.7, $F10.7 < 100$, $100 < F10.7 < 170$, and $F10.7$
66 > 170). Statistical averaging for each grid allowed the mean, median, and standard deviation for each
67 bin to be calculated whilst bi-linear interpolation allowed flux predictions to be made for any chosen
68 input values. The model also returned predictions of the 5th, 25th, 75th, and 95th percentiles of the
69 flux values for any chosen combination of input values. Hence, in the published version of the original
70 model, the user can input a particular energy, local time, and value of the Kp index, and the model will
71 return a prediction of the electron flux and the ion flux to be encountered at GEO for the chosen
72 energy, chosen local time, and chosen Kp index, at four different levels of solar activity [*Denton et al.*,
73 2015].

74

75 The bulk morphology of the electrons and ions at GEO, in the energy range sampled by LANL/MPA,

76 has previously been shown to be well-correlated with the level of magnetospheric convection [*Korth et*
77 *al.*, 1999; *Denton et al.*, 2005; 2007; *Lavraud et al.*, 2005]. And since the Kp index, with a 3 hour
78 cadence, is a very good proxy for this magnetospheric convection [*Thomsen*, 2004], then the original
79 model predicted fluxes that were in reasonably good agreement with observations. However, two
80 disadvantages arise from use of the Kp index, with particular regard to predictions. Firstly, the K
81 index (from which Kp is derived) is an Earth-based index, constructed from magnetometer
82 measurements of the horizontal component of the terrestrial magnetic field. Hence, estimates of flux at
83 Earth (geosynchronous orbit), based on the Kp index, are only available on an instantaneous basis (i.e.
84 a 'nowcast'), rather than being true advance predictions (i.e. a 'forecast'). Secondly, the fluxes at GEO
85 are regularly observed to fluctuate much more rapidly than three-hour time cadence of the Kp index,
86 typically in response to dynamic changes in the solar wind with timescales much less than one hour in
87 duration. Thus, our desire for a new and improved predictive model is driven by the following criteria:
88 (i) that the new model should be driven by some set of parameters that are regularly measured in the
89 solar wind, upstream of the Earth, and thus provide at least a one-hour time interval between prediction
90 of the fluxes and arrival of the fluxes at GEO, (ii) that the activity parameters should be capable of a
91 time cadence of at least one hour, and preferably as short as one minute, and (iii) that the new model
92 produce flux predictions that are, in the majority of cases, comparable with, or better than, the previous
93 version of the model. As outlined below, the results summarized in this study indicate that we have
94 largely achieved our intended aims by parameterizing the new model with the measured value of the
95 solar-wind electric field at the L1 point. The developmental methodology used in formulating the new
96 model is outlined in detail in Section 2, comparisons between model predictions and in-situ
97 observations of fluxes are made, along with goodness-of-fit calculations, in Section 3, and a discussion
98 of the strengths and weaknesses of the current model, and a summary, are provided in Section 4.

100 **2. Model Methodology**

101 The methodology followed in generating a new model of the fluxes at GEO is very similar to that used
 102 in the previous model, and described in detail in *Denton et al.* [2015]. The dataset for the model
 103 comprises ~82 satellite-years of electron and ion observations made between 1989 and 2007 by the
 104 LANL/MPA instruments flown on seven satellites at GEO. All flux measurements during this period
 105 are utilized when concurrent solar-wind measurements are available in the OMNI2 database [*King and*
 106 *Papitashvili, 2005*]. Periods when individual satellites are outside the magnetopause (usually during
 107 extremely high solar-wind pressure events) are excluded from the binning. One difference between
 108 this study and the previous *Denton et al.* [2015] study is that here we do not remove periods of
 109 exceptionally high spacecraft surface charging. The methodology to correct the particle energies
 110 resulting from the charging (due to acceleration towards the spacecraft, or repulsion away from the
 111 spacecraft), is considered robust [*Thomsen et al., 1999*].

112

113 In brief, all available flux values (for electrons and ions) are binned into one-hour width bins in local
 114 time, 40 logarithmically spaced bins in energy (from 1 eV to 40 keV), and 32 equal-width bins in -
 115 $v_{sw}B_z$ (from -8000 to +8000 $\mu\text{V m}^{-1}$). This binning yields a set of three-dimensional data-cubes that
 116 contain the mean, standard-deviation, and the 5th, 25th, 50th, 75th, and 95th percentiles of all data
 117 contributing to each bin (see Figure 3). In order to provide the average conditions at any local-time,
 118 energy or $-v_{sw}B_z$, bi-linear interpolation (with respect to the chosen local time and energy) and linear
 119 interpolation (with respect to $-v_{sw}B_z$) is used. The local time (in hours), the energy (in eV), and the
 120 negative product of the solar wind flow speed (v_{sw} in units of km s^{-1}) and the z-component of the solar-
 121 wind magnetic field (B_z in GSM coordinates in units of nT), are chosen by the user. This product is

122 the solar-wind electric field ($-v_{sw} \times B_z$) in units of $\mu\text{V m}^{-1}$ and in this parameterization, a solar wind
123 speed of 450 km s^{-1} in a magnetic field value of $B_z = -14.7 \text{ nT}$ (GSM) yields an electric field of $+6615$
124 $\mu\text{V m}^{-1}$. It is generally accepted that, as with the Kp index, the solar-wind electric field is reasonably
125 well-correlated with activity in the magnetosphere, although the detailed micro-physics that control the
126 coupling are beyond the scope of this paper (see *Newell et al.* [2007]; *Borovsky* [2013; 2014];
127 *McPherron et al.* [2015] for further discussions on this topic).

128
129 Although it is planned to evaluate more advanced coupling functions in future, use of the $-v_{sw} B_z$
130 parameter has the advantage of being well-known in the science and operational community, easily
131 computed, and widely available over the duration of the LANL/MPA dataset. In addition, this
132 parameter will be available in future via the continued operation of the ACE and DSCOVR satellites.
133 Solar wind data, propagated to the magnetopause, are taken from the high-resolution OMNI2 database
134 [*King and Papitashvili*, 2005] and MPA flux values are only included in the binning if solar-wind data
135 are available at the time of each data-point. To ensure a sufficient amount of data in each bin we have
136 limited the binning to $-v_{sw} B_z$ values between -8000 and $+8000 \mu\text{V m}^{-1}$, and do not provide separate
137 predictions for different values of the F10.7 index. [Note: The maximum flux variation between solar
138 maximum and solar minimum in the previous model was around a factor of 2, and only that large for a
139 small range of energies (\sim few keV). It is envisaged that users who will have a particular interest in
140 solar cycle effects will be able to examine the F10.7 variations in the previous model to gain insight
141 into the expected small changes with F10.7 in the new model].

142
143 The mean, standard-deviation, and the 5th, 25th, 50th, 75th, and 95th percentiles in each bin are calculated
144 for ions and electrons. Figure 2 contains a schematic representation of the binning process and Figure

145 3 shows the results of this binning for the mean electron flux, and the mean ion flux, at two example
 146 energies. The plots in this figure demonstrate how differences in the orientation of the interplanetary
 147 magnetic field direction (IMF), either northwards or southwards, radically change the average
 148 measured flux at GEO for both the electrons and the ions. Clearly, such differences are neglected
 149 when only considering the overall level of convection (proxied by Kp) as is the case in our previous
 150 model [Denton *et al.*, 2015]. Thus, we expect an increase in the prediction accuracy of the new model
 151 as a result. Figure 4 contains example surface plots showing the electron and ion flux variability, as a
 152 function of energy and local-time, for the one particular case when $-v_{sw}B_z = -2000 \mu\text{V m}^{-1}$.

153

154 3. Comparison of Model Predictions with Observations

155 In comparing observations with model predictions the aim is to evaluate both the general level of
 156 prediction ability of the model (goodness-of-fit), and also the incident solar-wind conditions for which
 157 the model predictions may be more, or less, accurate. Here, model predictions are compared against
 158 two different sets of observations - those provided by the LANL/MPA instruments themselves and
 159 those from the independent CEASE-II instrument [Dichter *et al.*, 1998] onboard the AMC-12 satellite,
 160 also located at GEO. The root-mean-squared deviation (RMSD) and the normalized root-mean-
 161 squared deviation (NRMSD) between the measured fluxes and the model predictions are calculated via
 162 the equation

$$163 \quad \text{NRMSD} = \text{RMSD} / (\bar{x}) = \sqrt{\frac{\sum_{x=1}^n [(x_{i, \text{model}} - x_{i, \text{measured}})^2]}{n}} / (\bar{x}) \quad (1)$$

164 where n is the number of data points over the range of the comparison and \bar{x} is the mean value of x
 165 over this range. Both NRMSD and RMSD are calculated in order to provide metrics with which to
 166 quantify the model accuracy (cf. Legates and McCabe Jr., 1999; Ganushkina *et al.*, 2015) although a

167 wide spectrum of other metrics may be used when comparing models to data [Koh *et al.*, 2012], each
168 with particular strengths and weaknesses. There are no universally accepted metrics for what
169 represents a 'good' NRMSD value, and certainly the calculated NRMSD values depend heavily on the
170 interval being studied. However, small values represent a better match between observations and
171 predictions than large values. The special case of RMSD=0 (NRMSD=0) would represent a perfect
172 forecast of the variation in the time-series being evaluated.

173

174 **3.1 Comparison with LANL/MPA observations**

175 A comparison of model predictions with the LANL/MPA observations at GEO is made for a calm
176 five-day period during 2004. Figure 5 contains electron observations and ion observations (at energies
177 ~32 keV) from the LANL-02A satellite (solid black line) along with model predictions from the Kp
178 version of the model (left column) and the new model driven by the solar wind electric field, $-v_{sw}B_z$
179 (right column). Although the original aim was to provide a model with a much higher temporal
180 resolution than the Kp model, on implementation it was found that rapid fluctuations in the $-v_{sw}B_z$
181 parameter resulted in rapid oscillations in model predictions. These do not accurately represent the
182 actual observations at GEO. Although it is unclear on what timescale the bulk magnetosphere
183 responds to changing solar-wind electric fields (likely a complicated function of particle energy,
184 species, time-history of the system, etc.), here the model results are smoothed with a five-minute box-
185 car average (this can be changed as required by the user) so as to smooth the oscillations in one-minute
186 high-resolution OMNI model input data. Note: the Kp model is naturally smoothed due to the 3-hour
187 cadence of the Kp index. The solid red line in Figure 5 is the predicted mean flux from the model, and
188 the solid purple line is the median. The 5th, 25th, 75th and 95th percentiles are indicated by the dashed
189 and dotted purple lines (the standard deviation is not shown). The Kp index and the $-v_{sw}B_z$ parameter

190 are plotted in the bottom row. The RMSD and NRMSD values for the model-data comparisons are
191 also provided in the top right of each plot. Both versions of the model provide a reasonably good fit to
192 the data with NRMSD values between ~ 0.14 - 0.25 . Little difference is apparent between the models
193 during these calm conditions, with the observed flux almost always falling within the 5th-to-95th
194 percentile range of the model predictions.

195

196 The plots shown in Figure 6 follow the same format but this time for ions and electrons with energies
197 ~ 10 keV) during a highly dynamic and disturbed period, also in 2004. The model predictions closely
198 follow the trend of the observations and for this period, even during some of the most dynamic changes
199 in the Kp index, and in $-v_{sw}B_z$. The NRMSD values are between ~ 0.15 - 0.21 for the new $-v_{sw}B_z$ model
200 and the original Kp model at these times, and these values are typical of a range of other energies
201 between ~ 1 - 40000 eV. Of course the $-v_{sw}B_z$ model also has the distinct advantage that it can make
202 flux-predictions ~ 1 hour prior to the event, provided the upstream solar-wind electric field value is
203 known. Again, the observed fluxes fall within the 5th-to-95th percentile range predicted by both
204 models during almost the entire period under study, although the sharp drop in the ion flux at the start
205 of day 94 is not predicted by either model.

206

207 A complete comparison of flux-predictions from the two different models at all observed energies can
208 be made examining energy-time spectrograms of flux values from in-situ observations, along with
209 simulated spectra from model predictions. Figure 7 contains electron (left column) and ion (right
210 column) flux spectrograms from the LANL-02A satellite (top row) with simulated spectra from the Kp
211 model (middle row) and the new $-v_{sw}B_z$ model (bottom row). The model spectra show the mean-flux
212 predictions from each model (although it is straight-forward to also evaluate the 5th, 25th 50th, 75th or

213 95th percentiles flux-predictions, along with the standard deviation). The orange dashed line indicates
214 local noon and the black dashed line indicates local midnight. Note: the observations of electron fluxes
215 below ~ 100 eV should be treated with caution due to the possible presence of photoelectrons and
216 secondary electrons contaminating the observations at these energies (see Fig. 7 top left panel).

217
218 Both the $-v_{sw}B_z$ and the Kp model flux spectra show many similarities to the observed LANL-02A
219 spectra with the model flux values being broadly comparable to the observations. However, it is clear
220 that there are significant differences at most energies. With respect to the electrons, the high fluxes
221 observed at energies up to 10 keV by LANL-02A after ~ 15 UT are not fully captured by either model.
222 The Kp model reproduces elevated fluxes at this time but their spatial structure is clearly affected by
223 the three-hour cadence of the Kp index. The $-v_{sw}B_z$ model reproduces rapid fluctuations in the fluxes
224 that result from changes in the solar-wind electric field, but in general these are not seen in the
225 observations prior to 15 UT. With respect to the ions, both the Kp and the $-v_{sw}B_z$ model reproduce the
226 appearance of a low-energy population (the ion plasmasphere) observed by LANL-02A but the model
227 fluxes are somewhat lower than actually observed.

228 229 **3.2 Comparison with independent AMC-12/CEASE-II observations**

230 As in evaluation of the previous model, in order to provide independent testing of the model veracity
231 (at least for the electron observations) we carry out a comparison of model predictions with electron
232 flux observations from the CEASE-II sensor onboard the AMC-12 satellite at GEO. This comparison
233 is made for DOY 180 during 2013, a particularly disturbed period during a geomagnetic storm where
234 Kp reached a maximum ~ 6 and the Dst index reached a maximum excursion ~ -100 at the start of the
235 day. As was previously noted [*Denton et al.*, 2015] there is a semi-constant offset between CEASE-II

236 fluxes and MPA model predictions and hence to account for this difference we multiply the CEASE-II
237 fluxes by a factor of ~ 15 at all times. Since no cross-calibration between the MPA and CEASE-II
238 instruments took place prior to launch, this adjustment is akin to on-orbit cross-calibration of the
239 fluxes. Note: for future comparison of model fluxes with measured fluxes from different satellites it
240 would be necessary to evaluate the need for use of an appropriate cross-calibration factor.

241

242 Figure 8 shows electron fluxes measured by the CEASE-II instrument during a 24-hour period in
243 color-spectrogram format, as a function of energy and time (top plot), along with the model electron
244 flux predictions from the Kp model (middle panel) and the new $-v_{sw}B_z$ model (bottom panel). The Kp
245 index and $-v_{sw}B_z$ are also shown, demonstrating the activity levels during this day. The predictions
246 from each model demonstrate that the broad features observed at GEO by out-of-sample instruments
247 such as CEASE-II, can be predicted, even during highly disturbed periods. The advantage of the $-$
248 $v_{sw}B_z$ model is that the fluxes to be encountered by the AMC-12 satellite can be predicted ~ 1 hour in
249 advance, given upstream solar wind measurements of the speed and z-component of the magnetic field.

250

251 **3.3 Spacecraft surface charging on LANL/MPA**

252 Along with the electron and ion fluxes, the MPA instruments also measure the electrostatic surface
253 potential on the LANL spacecraft, relative to the ambient plasma [Thomsen *et al.*, 1999]. Depending
254 on their individual design and construction details, spacecraft can charge positive or negative
255 [DeForest, 1972; Garrett, 1981; Farthing *et al.*, 1982; Lanzerotti *et al.* 1998; Thomsen *et al.* 2013]. In
256 the case of the LANL satellites, the greatest level of charging occurs during hours of eclipse when the
257 surface potential can reach 1000s Volts (negative) with respect to the ambient plasma. Such elevated
258 charging can be detected by the observation of an ion-line in the ion flux measurements. This occurs

259 due to positive particles that are accelerated towards the spacecraft by the negative charge on the
260 spacecraft. An example of such an ion-line can be seen between ~18-22 UT in the ion flux
261 observations in Figure 7.

262

263 Since the MPA instruments regularly measure the spacecraft charging, it is straight-forward to extend
264 the current flux model to include predictions of spacecraft charging, via similar methodology as that
265 used for the fluxes. Although each spacecraft charges differently depending on its construction, the
266 environmental conditions that give rise to dangerous levels of surface charging on one satellite are
267 likely to pose a danger to other satellites passing through the region. Figure 9 shows the mean
268 measured (negative) surface potential from all seven LANL satellites, in the same format as that used
269 to calculate the flux distributions (as a function of $-v_{sw}B_z$) shown in Figure 3. It is clear from Figure 9
270 that the most severe surface charging of the LANL satellites occurs during southwards IMF- B_z and at
271 spatial locations around local midnight. Charging is greatly reduced during periods of positive IMF-
272 B_z . (cf. plots of spacecraft charging from LANL/MPA as a function of K_p , Dst , and v_{sw} given in
273 *Denton and Borovsky [2012]*). The model predicts the level of surface charging on the LANL satellites
274 by carrying out a bi-linear interpolation between the mean surface charging levels in the appropriate
275 bins in Figure 9, based on the prevailing solar-wind conditions and the satellite local time. In this
276 respect the model predictions of surface charging are calculated similarly to the model fluxes. It is
277 planned that this predictive capability of model will be further developed in future versions.

278

279 **4. Discussion and Summary**

280 The ultimate goal of much "space weather" research is to *accurately* predict the conditions to be
281 encountered by orbital hardware systems as far in advance as possible. Of course, it is nigh-on

282 impossible for 100% accurate predictions to ever be achieved. However, by carrying out the work
283 outlined above, we aimed to achieve quantitative predictions that allow hardware operators and
284 scientific modelers the ability to predict fluxes in advance given knowledge of upstream solar wind
285 parameters. The absolute flux values (Figures 5 and 6), and the flux spectra (Figures 7 and 8), show
286 that the $-v_{sw}B_z$ model provides reasonably accurate flux predictions at GEO ~1 hour in advance,
287 providing knowledge of the solar wind electric field (e.g. from the ACE or DSCOVR satellites) is
288 available. Such knowledge is available in real-time (e.g. from the Space Weather Prediction Center
289 (<http://www.swpc.noaa.gov/>)).

290

291 With a view to potential changes that could improve future flux forecasts, it is important to be guided
292 by knowledge of the physics of the inner magnetosphere. The current model, and the previous version,
293 considered neither the time-history of the magnetosphere at the time of the predictions, or the explicit
294 transport times for plasma to migrate from the solar wind to the various locations around GEO. Drift
295 times are energy dependent, and also dependent on the local convection strength. Such potentially
296 non-linear effects can be estimated but are not known without complex particle tracing calculations.
297 *Denton and Borovsky* [2009] estimated transport timescales from the solar wind to various locations
298 around GEO with timescales being of the order of 0 h to 17.5 h. In addition, *Lavraud et al.* [2006]
299 demonstrated the importance of the time-history of the system with respect to plasma conditions at
300 GEO by examining the build up of cold, dense plasma during extended periods of northwards IMF.
301 Our aim is to explore inclusion of both of these effects in future versions of the model.

302

303 With regard to operational uses of the model, *Thomsen et al.* [2013] demonstrated that satellite surface
304 charging is strongly correlated with periods when the electron flux at energies between 5-10 keV

305 exceeds a particular threshold. That study found that satellite surface charging was most likely to
306 occur during intervals when the electron flux at 8 keV exceeded a flux threshold of $1.4 \times 10^3 \text{ cm}^{-2} \text{ s}^{-1}$
307 $\text{str}^{-1} \text{ eV}^{-1}$. Armed with this knowledge, one possible use of the model would be to: (i) determine
308 appropriate cross-calibration factors between the model (based on MPA) and fluxes measured by the
309 chosen satellite; (ii) use upstream values of solar-wind electric field to search for intervals when the
310 predicted electron fluxes at 8 keV exceeded a flux threshold; (iii) expect elevated surface charging to
311 be more likely during such intervals.

312

313 The model provides good agreement with in-sample MPA observations and (with appropriate on-orbit
314 cross-calibration) with independent out-of-sample observations from the CEASE-II detector onboard
315 AMC-12. It is hoped that the model will prove useful to the community of orbital hardware designers
316 and satellite operators, as well as to the scientific community who use fluxes at GEO as inputs to
317 physical models.

318

319 In summary:

320

321 1. A new model of the electron fluxes and ion fluxes at GEO, which uses the solar-wind electric field
322 as input, has been developed. The model provides a ~1 hour advanced forecast of the fluxes at GEO in
323 the energy range ~1 eV to ~ 40 keV.

324

325 2. The model provides forecasts of the fluxes at GEO that are comparable in accuracy to the previous
326 model, driven by the Kp index. The main benefit from the new model is the ability to predict the
327 fluxes in advance.

328

329 3. The model results are robust, during both quiet times and highly disturbed storm-times, as measured
330 by the Root-Mean-Squared-Deviation (RMSD) and the Normalized Root-Mean-Squared Deviation
331 (NRMSD). Observed fluxes are found to almost always fall between the 5th and 95th percentiles of
332 the model predictions.

333

334 4. A new forecasting capability for spacecraft surface charging on the LANL satellites is introduced.
335 Further development of this capability is planned for the future.

336

337 The model is freely available to users under the GNU General Public License v3.0 by contacting the
338 author directly or via the model webpage at <http://gemelli.space-science.org/mdenton/>.

339

340 **Acknowledgements**

341 The authors gratefully acknowledge the OMNI database for the solar wind and geophysical parameters
342 used in this study. We also acknowledge use of the list of satellites at GEO maintained by Eric
343 Johnston at www.satsig.net. This work was partially supported by the Space Hazards Induced near
344 Earth by Large, Dynamic Storms (SHIELDS) project, funded by the U.S. Department of Energy
345 through the LANL/LDRD Program under contract DE-AC52-06NA25396. MHD wishes to thank J.
346 Denton for help with the Kp version of the model and comments on the current manuscript. MPA data
347 are available by contacting the PI, Mike Henderson, at mghenderson@lanl.gov. The model, written in
348 FORTRAN, is available from MHD at mdenton@space-science.org, or via download at
349 <http://gemelli.space-science.org/mdenton/>.

350 **References**

- 351 Bame, S. J., D. J. McComas, M. F. Thomsen, B. L. Barraclough, R. C. Elphic, J. P. Glore, J. C.
352 Chavez, E. P. Evans and F. J. Wymer, *Rev. Sci. Instrum.*, 64, 1026-1033, 1993.
- 353 Bartels, J., N. A. H. Heck, and H. F. Johnstone, The three-hour-range index measuring geomagnetic
354 activity, *J. Geophys. Res.*, 44, 411-454, 1939.
- 355 Borovsky, J. E., Physical improvements to the solar wind reconnection control function for the Earth's
356 magnetosphere, *J. Geophys. Res. Space Physics*, 118, 2113–2121, 2013.
- 357 Borovsky, J. E., Canonical correlation analysis of the combined solar wind and geomagnetic index data
358 sets, *J. Geophys. Res. Space Physics*, 119, 5364–5381, 2014.
- 359 DeForest, S. E., Spacecraft charging at synchronous orbit, *J. Geophys. Res.*, 77, 651-659, 1972.
- 360 Denton, M. H., M. F. Thomsen, V. K. Jordanova, M. G. Henderson, J. E. Borovsky, J. S. Denton, D.
361 Pitchford, and D. P. Hartley, An empirical model of electron and ion fluxes derived from
362 observations at geosynchronous orbit, *Space Weather*, 13, 2015.
- 363 Denton, M. H., and J. E. Borovsky, Magnetosphere response to high-speed solar-wind streams: A
364 comparison of weak and strong driving and the importance of extended periods of fast solar wind,
365 *J. Geophys. Res.*, 117, A00L05, doi:10.1029/2011JA017124, 2012.
- 366 Denton, M. H., and J. E. Borovksy, The superdense plasma sheet in the magnetosphere during high-
367 speed-stream-driven storms: Plasma transport timescales *J. Atmos. Sol-Terr. Phys*, 71, 1045-1058,
368 2009
- 369 Denton, M. H., M. F. Thomsen, B. Lavraud, M. G. Henderson, R. M. Skoug, H. O. Funsten, J.-M.
370 Jahn, C. J. Pollock, and J. M. Weygand, Transport of plasma sheet material to the inner
371 magnetosphere, *Geophys. Res. Lett.*, 34, L04105, doi:10.1029/2006GL027886, 2007.
- 372 Denton, M. H., M. F. Thomsen, H. Korth, S. Lynch, J. C. Zhang and M. W. Liemohn, Bulk plasma

373 properties at geosynchronous orbit, *J. Geophys. Res.*, 110, A07223, 2005.

374 Dichter, B. K., J. O. McGarity, M. R. Oberhardt, V. T. Jordanov, D. J. Sperry, A. C. Huber, J. A.
375 Pantazis, E. G. Mullen, G. Ginet, and M. S. Gussenhoven, Compact Environmental Anomaly
376 Sensor (CEASE): A novel spacecraft instrument for in situ measurements of environmental
377 conditions, *IEEE Trans. Nucl. Sci.*, 45, 2758–2764, 1998.

378 Farthing, W. H., J. P. Brown, and W. C. Bryant, Differential spacecraft charging on the geostationary
379 operational satellites, NASA Tech. Memo, NASA TM-83908, 1982.

380 Ganushkina, N. Y., O. A. Amariutei, D. Welling, and D. Heynderickx, Nowcast model for low-energy
381 electrons in the inner magnetosphere, *Space Weather*, 13, 16–34, doi:10.1002/2014SW001098,
382 2015.

383 Ganushkina, N. Y., M. W. Liemohn, O. A. Amariutei, and D. Pitchford, Low-energy electrons (5–50
384 keV) in the inner magnetosphere, *J. Geophys. Res. Space Physics*, 119, 246–259,
385 doi:10.1002/2013JA019304, 2014.

386 Ganushkina, N. Y., O. Amariutei, Y. Y. Shpritz, and M. Liemohn, Transport of the plasma sheet
387 electrons to the geostationary distances, *J. Geophys. Res. Space Physics*, 118, 82–98,
388 doi:10.1029/2012JA017923, 2013.

389 Garrett, H. B., The charging of spacecraft surfaces, *Rev. Geophys.*, 19, 577-616, 1981.

390 Ginet, G. P., T. P. O’Brien, S. L. Huston, W. R. Johnston, T. B. Guild, R. Friedel, C. D. Lindstrom, C.
391 J. Roth, P. Whelan, R. A. Quinn, D. Madden, S. Morley, and Yi-Jiun Su, AE9, AP9 and SPM:
392 New Models for Specifying the Trapped Energetic Particle and Space Plasma Environment, in *The*
393 *Van Allen Probes mission*, eds N. Fox and J. L. Burch, Springer, doi:10.1007/978-1-4899-7433-4,
394 ISBN: 978-1-4899-7432-7, 2014.

395 Hartley, D. P., M. H. Denton, and J. V. Rodriguez, Electron number density, temperature, and energy

396 density at GEO and links to the solar wind: A simple predictive capability, *J. Geophys. Res. Space*
397 *Physics*, 119, 4556–4571, doi:10.1002/2014JA019779, 2014.

398 Jordanova, V. K., C. J. Farrugia, L. Janoo, J. M. Quinn, R. B. Torbert, K. W. Ogilvie, R. P. Lepping, J.
399 T. Steinberg, D. J. McComas, and R. D. Belian, October 1995 magnetic cloud and accompanying
400 storm activity: Ring current evolution, *J. Geophys. Res.*, 103, 79, 1998.

401 Jordanova, V. K., L. M. Kistler, M. F. Thomsen, and C. G. Mouikis, Effects of plasma sheet variability
402 on the fast initial ring current decay, *Geophys. Res. Lett.*, 30(6), 1311,
403 doi:10.1029/2002GL016576, 2003.

404 Katus, R. M., M. W. Liemohn, E. L. Ionides, R. Ilie, D. Welling, and L. K. Sarno-Smith, Statistical
405 analysis of the geomagnetic response to different solar wind drivers and the dependence on storm
406 intensity, *J. Geophys. Res.*, doi:10.1002/2014JA020712, 2014.

407 King, J. H., and N. E. Papitashvili, Solar wind spatial scales in and comparisons of hourly Wind and
408 ACE plasma and magnetic field data, *J. Geophys. Res.*, 110, A02104, 2005.

409 Koh, T.-Y. S. Wang, and B. C. Bhatt, A diagnostic suite to assess NWP performance, *J. Geophys.*
410 *Res.*, 2012, 117, D13109, 2012.

411 Korth, H., M. F. Thomsen, J. E. Borovsky, and D. J. McComas, Plasma sheet access to
412 geosynchronous orbit, *J. Geophys. Res.*, 104, 25,047–25,061, 1999.

413 Lanzerotti, L. J., C. Breglia, D. W. Maurer, G. K. Johnson III, and C. G. MacLennan, Studies of
414 spacecraft charging on a geosynchronous telecommunications satellite, *Adv. Space Res.*, 22, 79-
415 82, 1998.

416 Lavraud, B., M. F. Thomsen, J. E. Borovsky, M. H. Denton, and T. I. Pulkkinen, Magnetosphere
417 preconditioning under northward IMF: Evidence from the study of coronal mass ejection and
418 corotating interaction region geoeffectiveness, *J. Geophys. Res.*, 111, A09208,

419 doi:10.1029/2005JA011566, 2006.

420 Lavraud, B., M. H. Denton, M. F. Thomsen, J. E. Borovsky, and R. H. W. Friedel, Superposed epoch
421 analysis of dense plasma access to geosynchronous orbit, *Ann. Geophys.*, 23, 2519–2529, 2005.

422 Legates, D. R., and G. J. McCabe Jr., Evaluating the use of “goodness-of-fit” Measures in hydrologic
423 and hydroclimatic model validation, *Water Resour. Res.*, 35(1), 233–241,
424 doi:10.1029/1998WR900018, 1999.

425 McPherron, R. L., T.-S. Hsu, and X. Chu, An optimum solar wind coupling function for the AL index.
426 *J. Geophys. Res. Space Physics*, 120, 2494–2515, doi: 10.1002/2014JA020619, 2015.

427 Newell, P. T., T. Sotirelis, K. Liou, C.-I. Meng, and F. J. Rich, A nearly universal solar wind-
428 magnetosphere coupling function inferred from 10 magnetospheric state variables, *J. Geophys.*
429 *Res.*, 112, A01206, doi:10.1029/2006JA012015, 2007.

430 O’Brien, T. P., J. E. Mazur, and J. F. Fennell, The Priority Mismatch Between Space Science and
431 Satellite Operations, *Space Weather*, 11, doi:10.1002/swe.20028, 2013

432 O’Brien, T. P., SEAES-GEO: A spacecraft environmental anomalies expert system for
433 geosynchronous orbit, *Space Weather*, 7, S09003, 2009.

434 O’Brien, T. P., and C. L. Lemon, Reanalysis of plasma measurements at geosynchronous orbit, *Space*
435 *Weather*, 5, S03007, doi:10.1029/2006SW000279, 2007.

436 Purvis, C. K., H. B. Garrett, A. C. Whittlesey, and N. J. Stevens, Design guidelines for assessing and
437 controlling spacecraft charging effects, *NASA Tech. Pap.* 2361, 1984.

438 Sicard-Piet, A., S. Bourdarie, D. Boscher, R. H. W. Friedel, M. Thomsen, T. Goka, H. Matsumoto, and
439 H. Koshiishi, A new international geostationary electron model: IGE-2006, from 1 keV to 5.2
440 MeV, *Space Weather*, 6, S07003, doi:10.1029/2007SW000368, 2008.

441 Thomsen, M. F., M. G. Henderson, and V. K. Jordanova, Statistical properties of the surface-charging

442 environment at geosynchronous orbit, *Space Weather*, 11, 237-244, 2013.

443 Thomsen, M. F., E. Noveroske, J. E. Borovsky, and D. J. McComas, Calculating the Moments from
444 Measurements by the Los Alamos Magnetospheric Plasma Analyzer, LA-13566-MS, Los Alamos
445 National Laboratory, 1999.

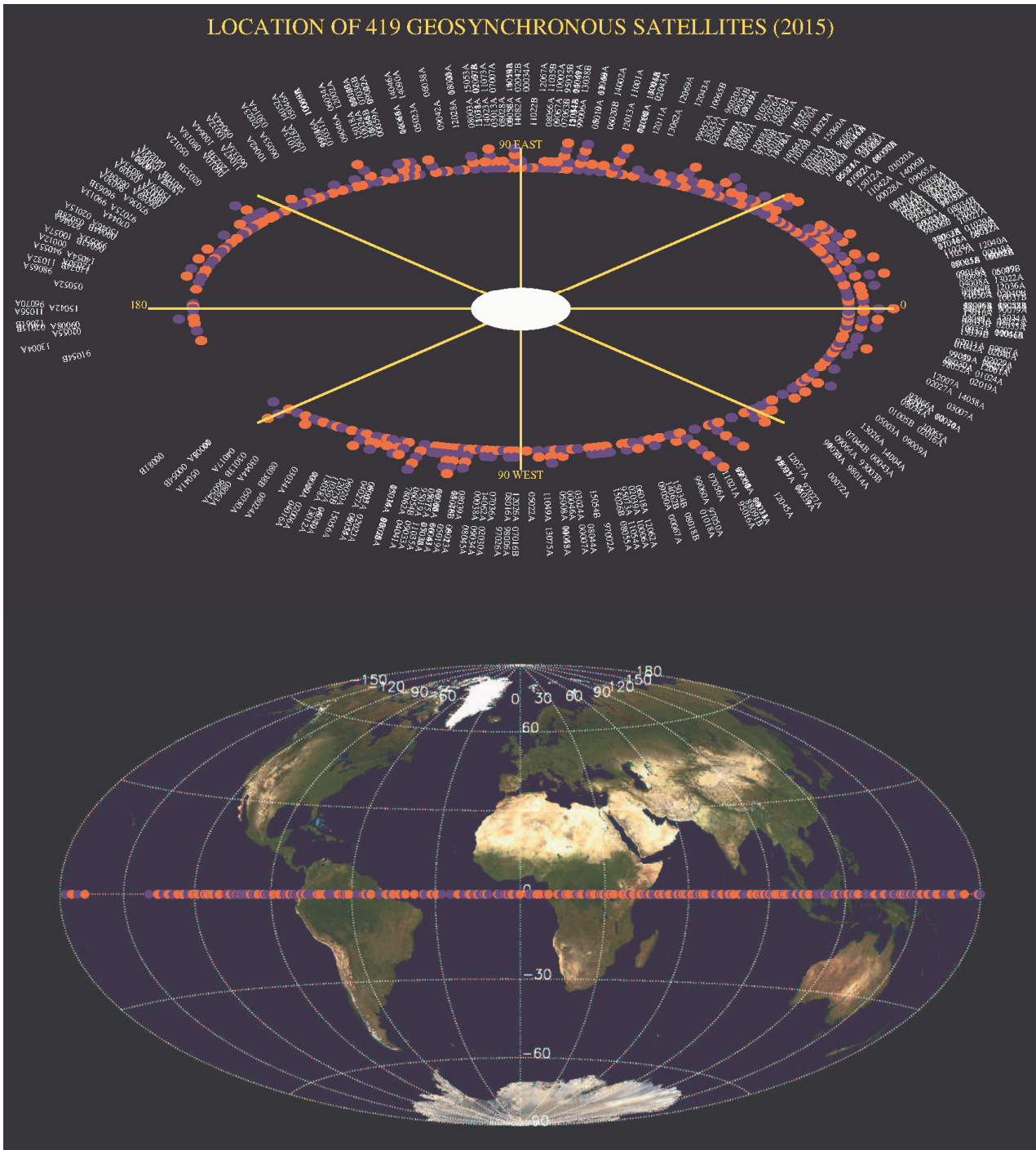
446 Thomsen, M. F., Why Kp is such a good measure of magnetospheric convection, *Space Weather*, 2,
447 S11004, doi:10.1029/2004SW000089, 2004.

448 Thomsen, M. F., M. H. Denton, B. Lavraud, and M. Bodeau, Statistics of plasma fluxes at
449 geosynchronous orbit over more than a full solar cycle, *Space Weather*, 5, S03004,
450 doi:10.1029/2006SW000257, 2007.

451 Zaharia, S., M. F. Thomsen, J. Birn, M. H. Denton, V. K. Jordanova, and C. Z. Cheng, Effect of storm-
452 time plasma pressure on the magnetic field in the inner magnetosphere, *Geophys. Res. Lett.*, 32,
453 L03102, doi:10.1029/2004GL021491, 2005.

454 Zaharia, S., V. K. Jordanova, M. F. Thomsen, and G. D. Reeves, Self-consistent modeling of magnetic
455 fields and plasmas in the inner magnetosphere: Application to a geomagnetic storm, *J. Geophys.*
456 *Res.*, 111, A11S14, doi:10.1029/2006JA011619, 2006.

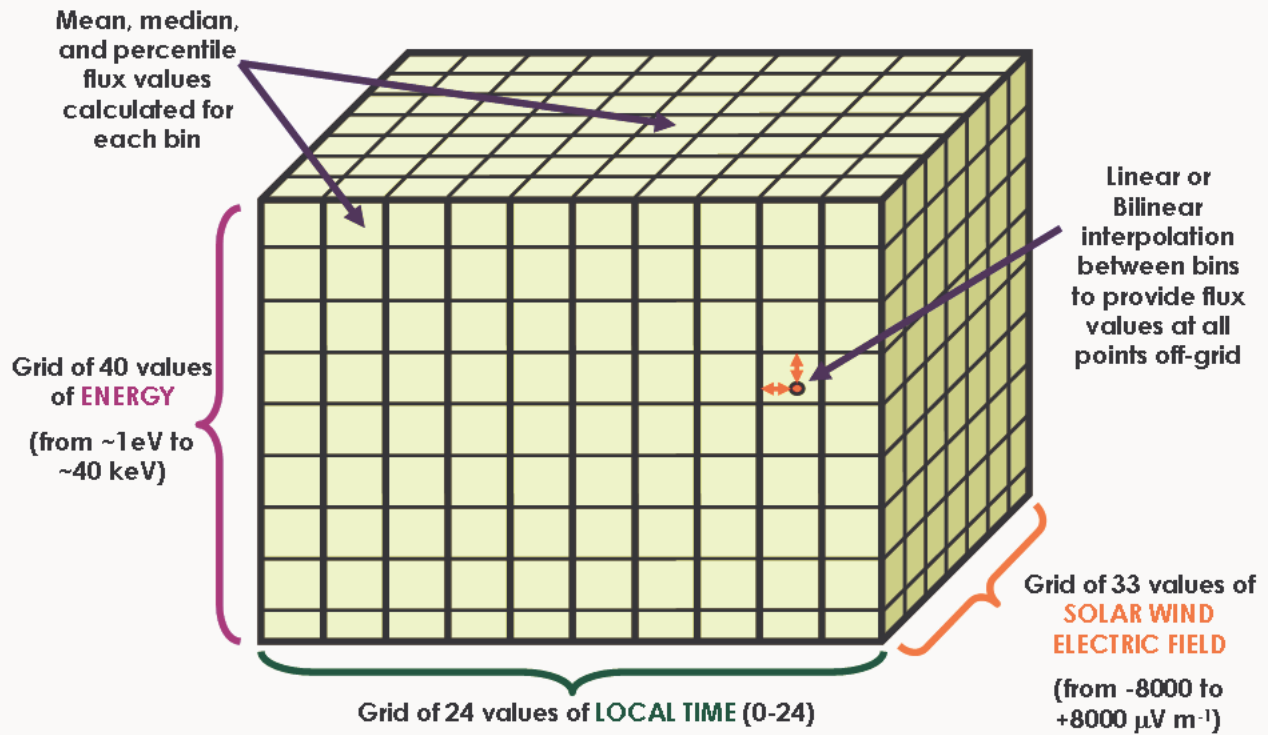
457



459
 460
 461
 462
 463
 464
 465
 466

Figure 1. (Top panel) The geographic (equatorial) locations of selected Earth-orbiting satellites located in a synchronous orbit east and west of the Greenwich meridian (0° longitude). Where more than one satellite is located at the same longitude (to 0.1 degree accuracy) the satellites are displayed radially outwards from GEO. (Bottom panel) The geographic equatorial footprint of the satellites on the Earth (Hammer-Aitoff projection). Note: no account is taken of the satellite inclination. Adjacent satellites are alternately displayed in red/blue for clarity.

Binning Scheme for $-vB_z$ Model



468

469

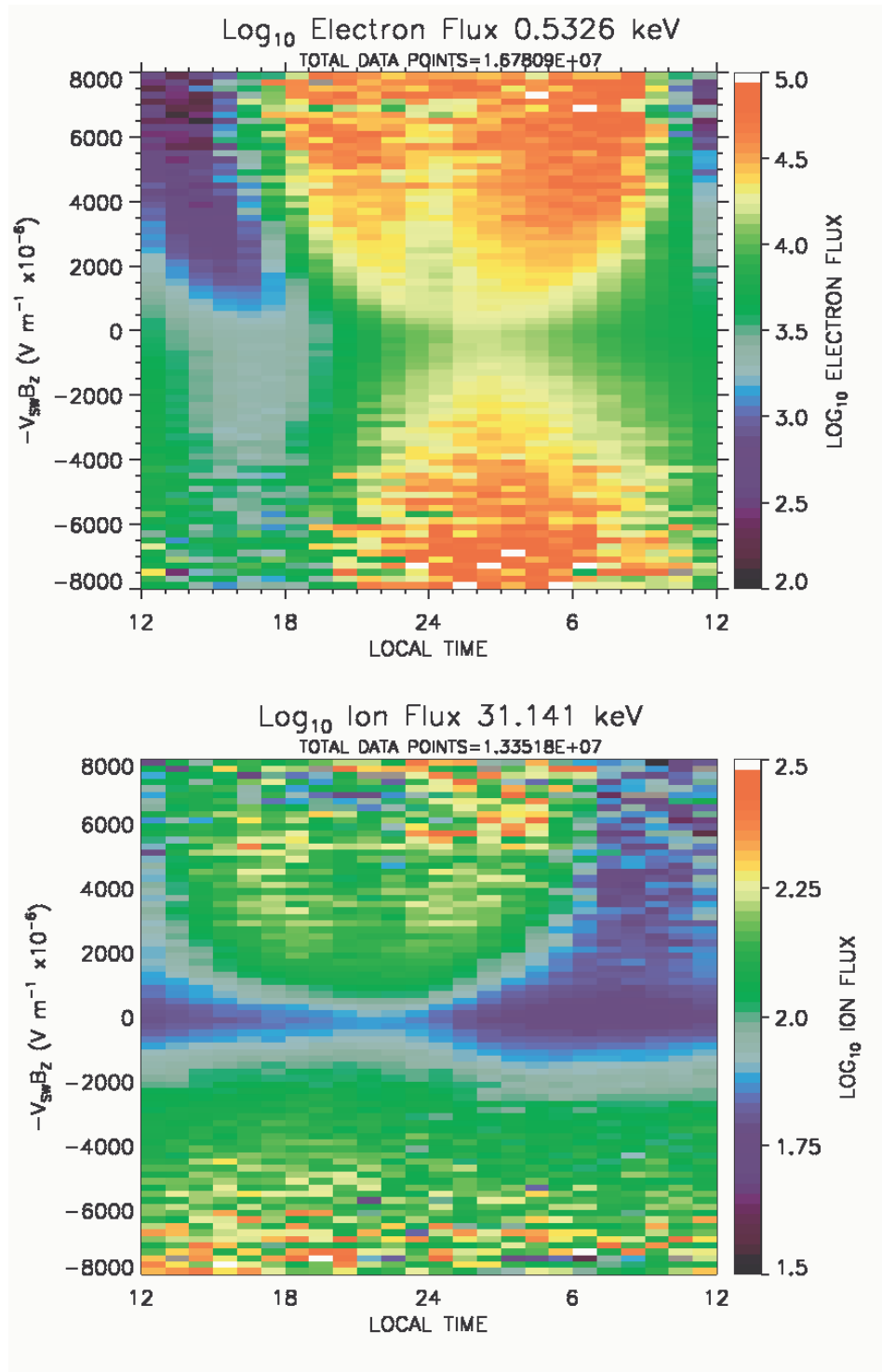
470

471

472

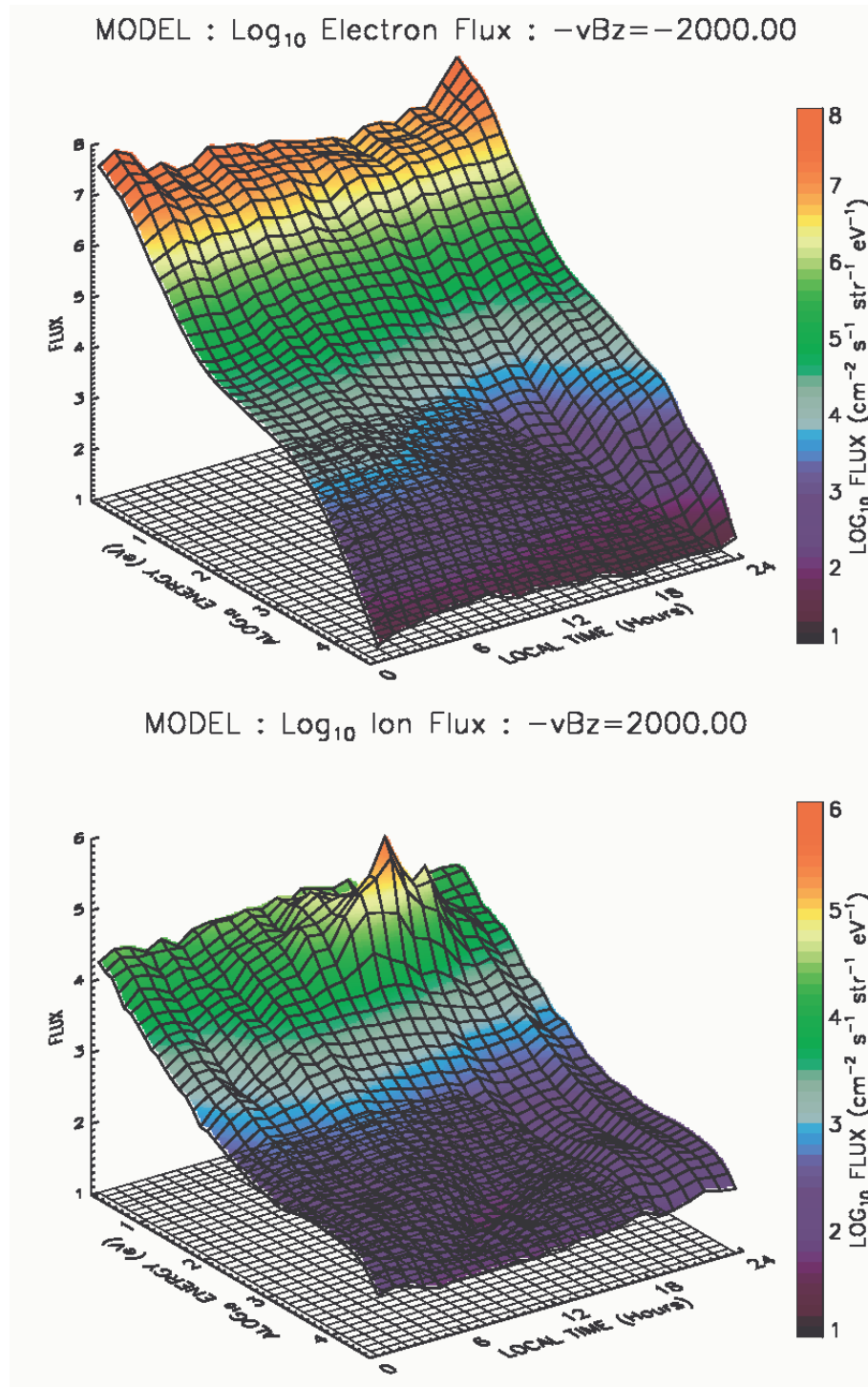
473

Figure 2. Schematic showing the binning scheme for the $-vB_z$ model. The three-dimensional model grid contains 40 energy bins (between 1 eV and 40 keV), 33 bins of $-vB_z$ values (-8000 to +8000), and 24 bins of local time (0-24), for both the ions and the electrons.



475
 476
 477
 478
 479
 480

Figure 3. Example mean fluxes from LANL/MPA as a function of solar-wind electric field and local time, for the electrons (top panel - 532.6 eV) and the ions (bottom panel - 31141 eV). These plots demonstrate the large difference in the average flux at GEO for cases where the IMF is northwards (negative $-v_{sw}B_z$) or southwards (positive $-v_{sw}B_z$).



482

483

484

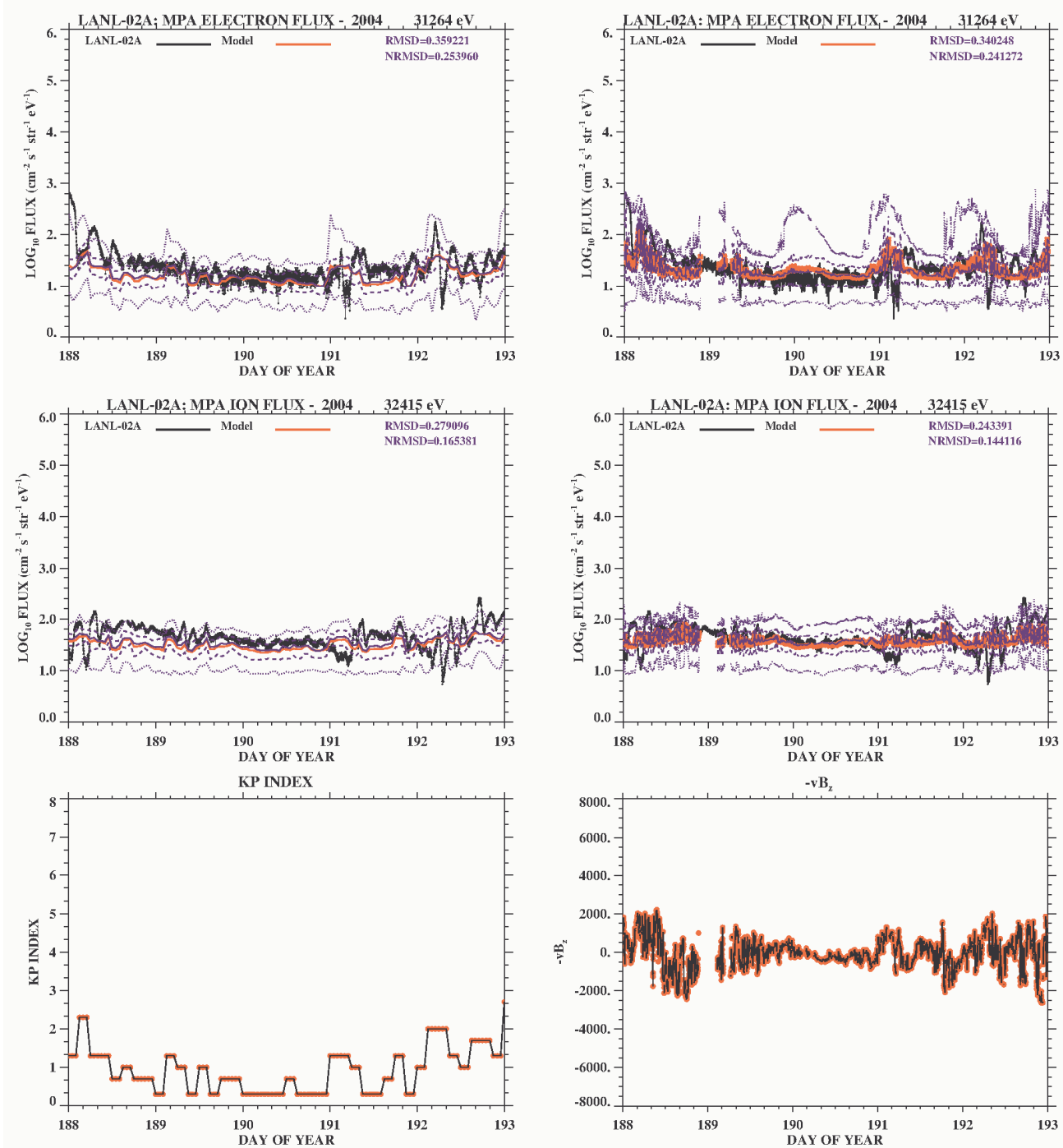
485

486

487

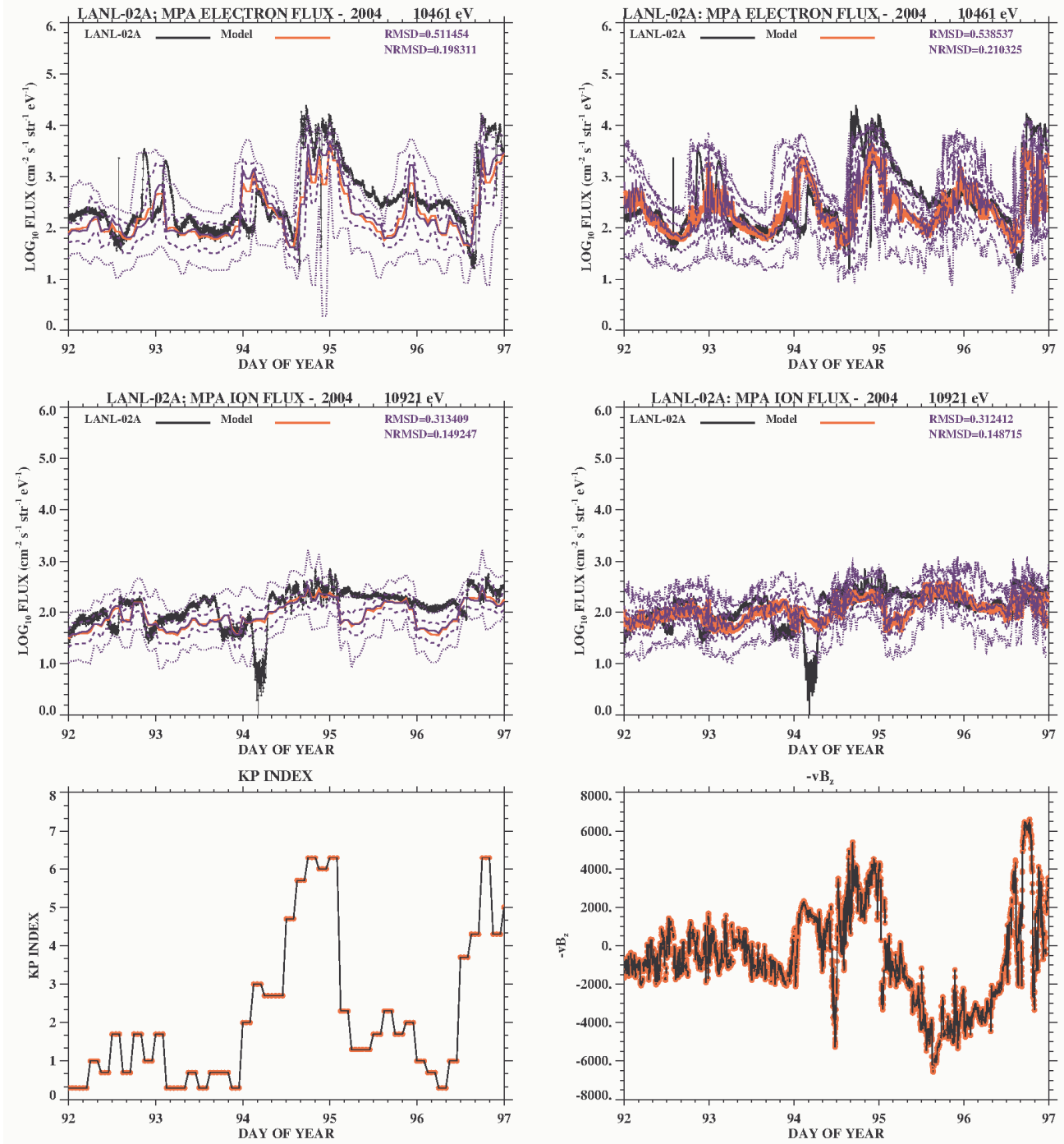
488

Figure 4. Example surfaces showing the model flux surfaces at two different values of $-v_{sw}B_z$, for electrons (top: $-vB_z = 2000 \mu\text{V m}^{-1}$.) and ions (bottom: $-vB_z = -2000 \mu\text{V m}^{-1}$.), as a function of energy and local time. Flux values at points off the grid can be computed via bi-linear interpolation between grid points, and subsequent linear interpolation between the discrete values of $-v_{sw}B_z$.



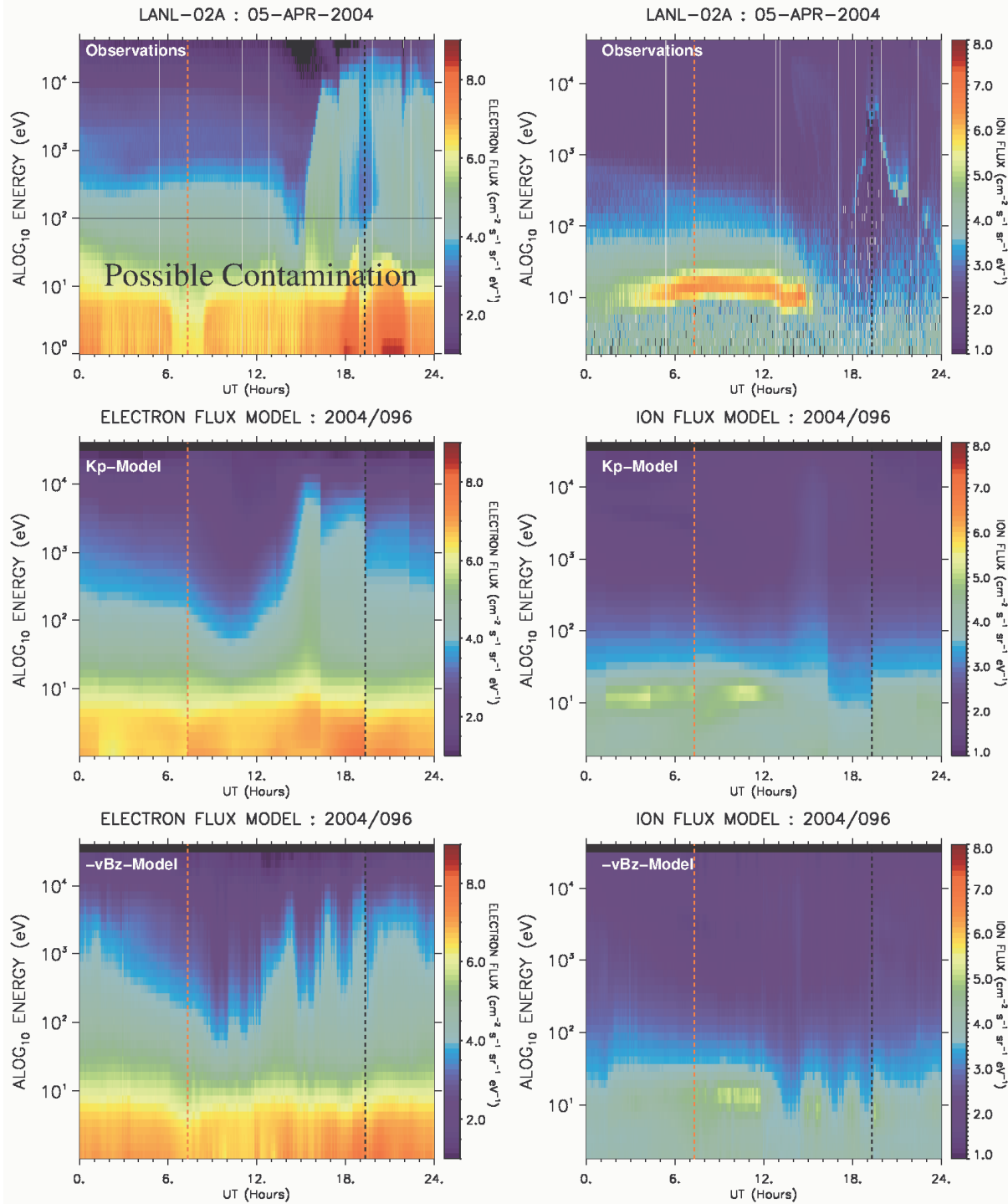
489
490

491 **Figure 5.** Example model results for five very calm days in 2004 for electrons and ions with energies
 492 ~ 32 keV. The left column shows the model predictions for electrons and ions using the Kp version of
 493 the model. The right column shows model predictions in the same format, but using the $-v_{sw}B_z$ version
 494 of the model. The black line shows the observations from the LANL-02A satellite. The solid red line
 495 is the predicted mean flux, and the solid purple line is the median. The 5th, 25th, 75th and 95th
 496 percentiles are indicated by the dashed and dotted purple lines. The Kp index and the $-v_{sw}B_z$ parameter
 497 are also shown in the bottom row.



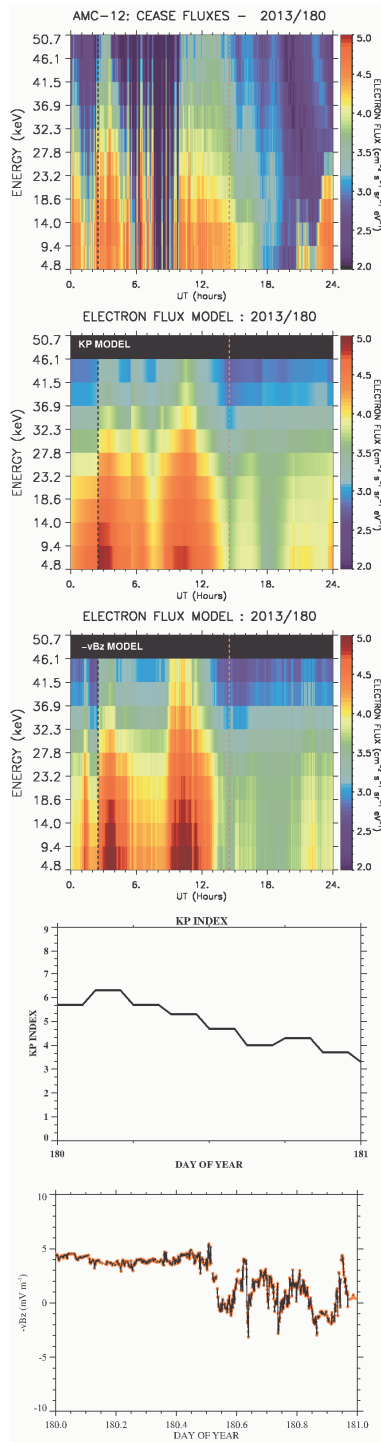
498
 499
 500
 501
 502
 503
 504
 505
 506

Figure 6. Example model results for five disturbed days in 2004 for electrons and ions with energies ~ 10 keV. The left column shows the model predictions for electrons and ions using the Kp version of the model. The right column shows model predictions in the same format, but using the $-v_{\text{sw}}B_z$ version of the model. The black line shows the observations from the LANL-02A satellite. The solid red line is the predicted mean flux, and the solid purple line is the median. The 5th, 25th, 75th and 95th percentiles are indicated by the dashed and dotted purple lines. The Kp index and the $-v_{\text{sw}}B_z$ parameter are also shown in the bottom row.



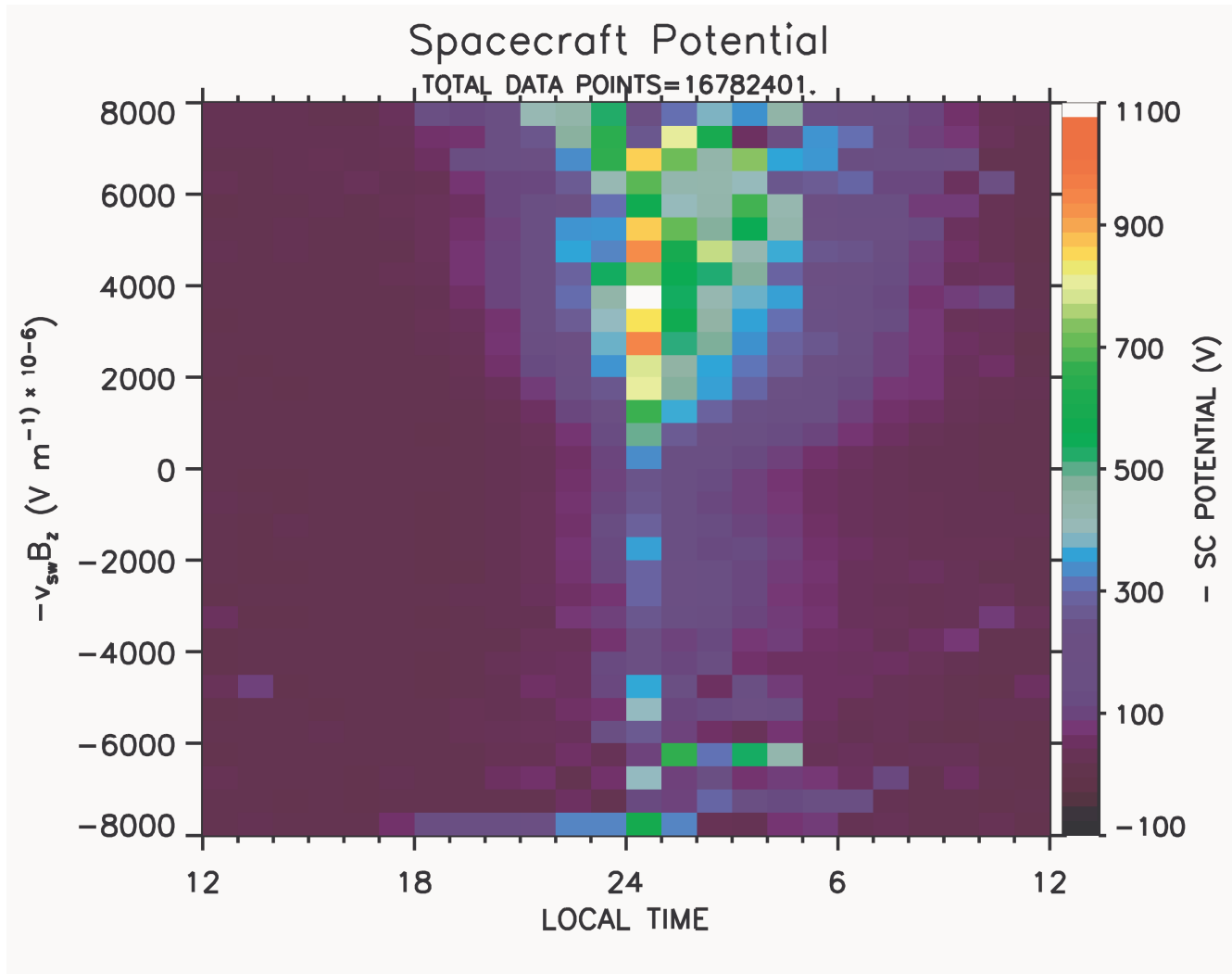
507
508
509
510
511
512
513

Figure 7. Comparison of particle flux observations (top row) with the Kp-model (middle row) and the new $-v_{sw}B_z$ model (bottom row) for electrons (left column) and ions (right column) on 5th April (day-of-year 96) in 2004. Large fluctuations occur in both the Kp index and in the $-v_{sw}B_z$ parameter on this day (see Figure 6). The orange line marks local noon and the black line marks local midnight in each plot.



514
515

516 **Figure 8.** Figure showing the CEASE-II electron flux observations from AMC-12 (top panel) on 29th
517 June (DOY- 180) in 2013 during disturbed geomagnetic activity. Also shown are the electron flux
518 predictions from the Kp model (middle panel), and the electron flux predictions from the $-v_{sw}B_z$ model
519 (bottom panel). Note: the CEASE-II electron fluxes have been multiplied by a constant factor (cross-
520 calibrated) to bring them into alignment with the LANL/MPA model fluxes. Kp and $-v_{sw}B_z$ are also
521 shown.



523

524

525 **Figure 9.** Showing the mean (negative) of the spacecraft potential measured by MPA spacecraft
 526 between 1990 and 2007 as a function of $-v_{sw}B_z$ and local time. The spacecraft surface potential is
 527 clearly most elevated around local midnight, and during southwards excursions of IMF- B_z .

528

529